

Recent advances in few-body nuclear reactions

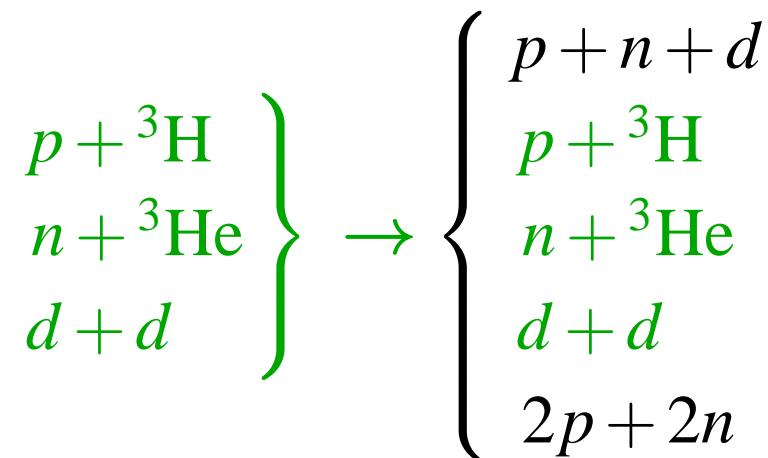
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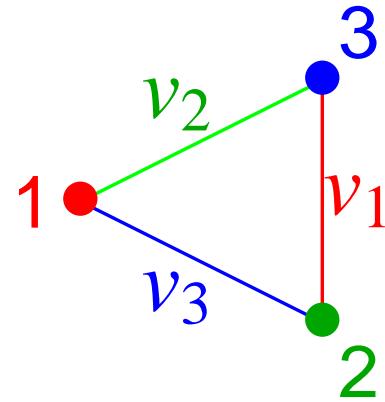
Outline

- core excitation:
extended Faddeev/AGS formalism
- 3-body nuclear reactions
 - $^{24}\text{Mg}(d, d')$
 - $^{10}\text{Be}(d, p)$, $^{11}\text{Be}(p, d)$, $^{11}\text{Be}(p, pn)$
 - $^{20}\text{O}(d, p)$
- 4-particle scattering
(p,n), (d,p) and (d,n) reactions in 4N system

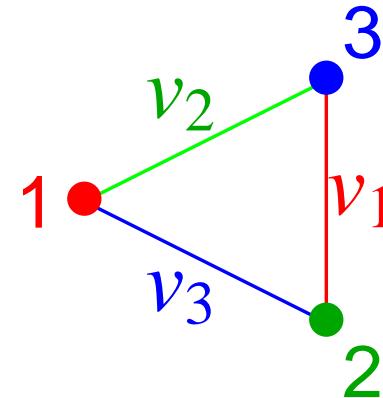


Three-particle system

Hamiltonian $H_0 + \sum_{\alpha} v_{\alpha}$



Three-particle system



Hamiltonian $H_0 + \sum_{\alpha} v_{\alpha}$

- Faddeev equations

$$(E - H_0 - v_{\alpha}) |\Psi_{\alpha}\rangle = v_{\alpha} \sum_{\sigma} \bar{\delta}_{\alpha\sigma} |\Psi_{\sigma}\rangle$$

$$|\Psi\rangle = \sum_{\alpha} |\Psi_{\alpha}\rangle$$

difficult to solve

exact solution of 3-body problem:

discrepancy with data

→ shortcomings of 3-body Hamiltonian

Alt, Grassberger, and Sandhas equations

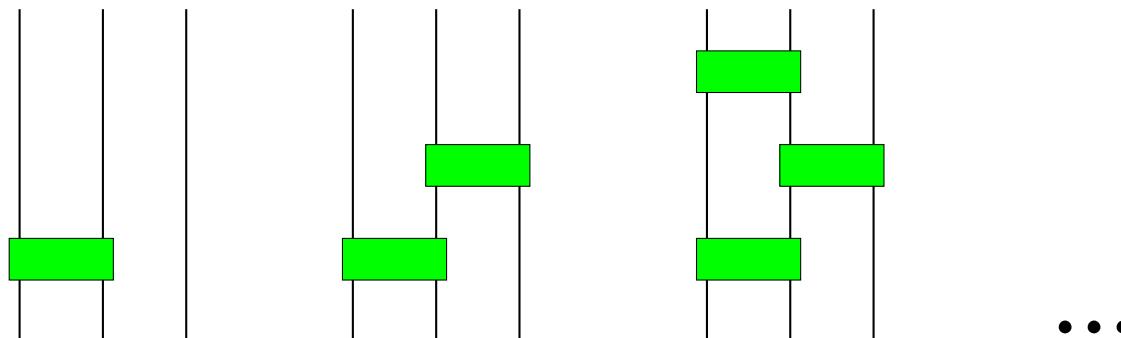
$$U_{\beta\alpha} = \bar{\delta}_{\beta\alpha} G_0^{-1} + \sum_{\sigma} \bar{\delta}_{\beta\sigma} T_{\sigma} G_0 U_{\sigma\alpha}$$

$$U_{0\alpha} = G_0^{-1} + \sum_{\sigma} T_{\sigma} G_0 U_{\sigma\alpha}$$

$$T_{\sigma} = v_{\sigma} + v_{\sigma} G_0 T_{\sigma}$$

$$G_0 = (E + i0 - H_0)^{-1}$$

channel states $(E - H_0 - v_{\alpha})|\phi_{\alpha}\rangle = 0$



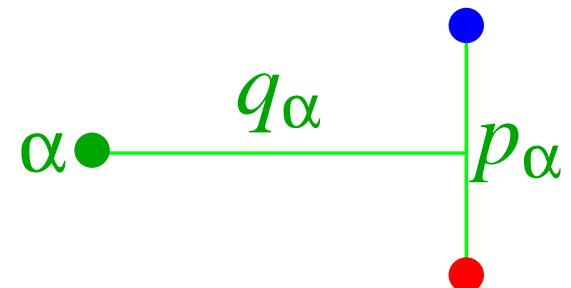
AGS equations with 3BF

$$V_{3BF} = \sum_{\alpha=1}^3 w_\alpha$$

$$\begin{aligned} U_{\beta\alpha} &= \bar{\delta}_{\beta\alpha} G_0^{-1} + \sum_{\gamma} \bar{\delta}_{\beta\gamma} T_{\gamma} G_0 U_{\gamma\alpha} \\ &\quad + w_{\alpha} + \sum_{\gamma} w_{\gamma} G_0 (1 + T_{\gamma} G_0) U_{\gamma\alpha} \end{aligned}$$

AGS equations: numerical solution

$$U_{\beta\alpha} = \bar{\delta}_{\beta\alpha} G_0^{-1} + \sum_{\sigma} \bar{\delta}_{\beta\sigma} T_{\sigma} G_0 U_{\sigma\alpha}$$

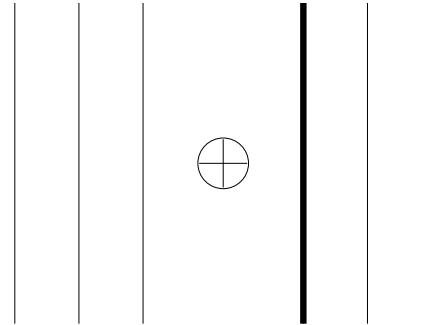


- 3 sets of Jacobi momenta
- momentum-space partial wave basis
- set of coupled 2-variable integral equations
- integrable singularities in kernel
- Coulomb interaction: screening and renormalization

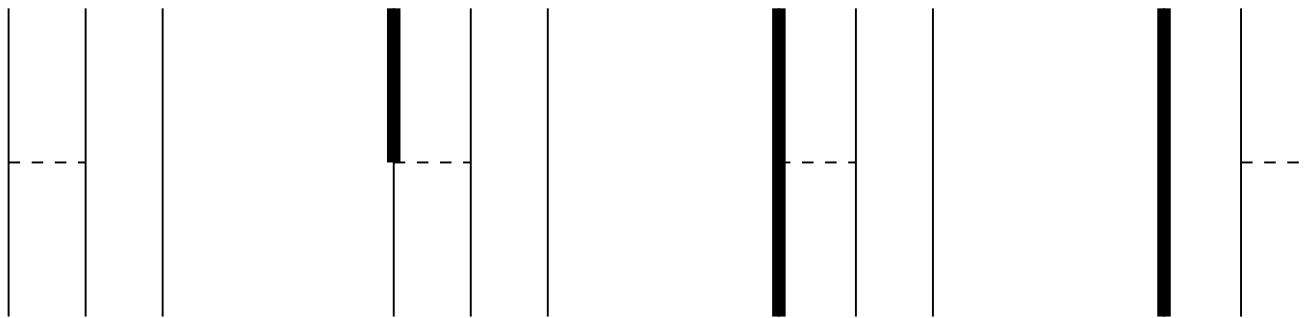
[PRC 71, 054005; PRC 72, 054004; PRC 74, 064001]

Core excitation (CX): extended Hilbert space

$$\mathcal{H} = \mathcal{H}_g \oplus \mathcal{H}_x$$

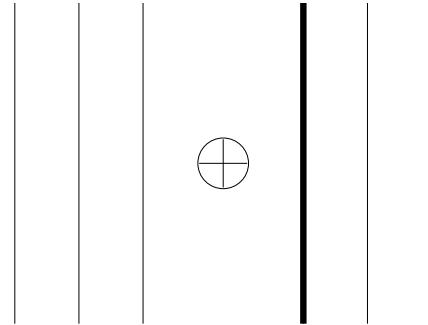


sector coupling by interaction

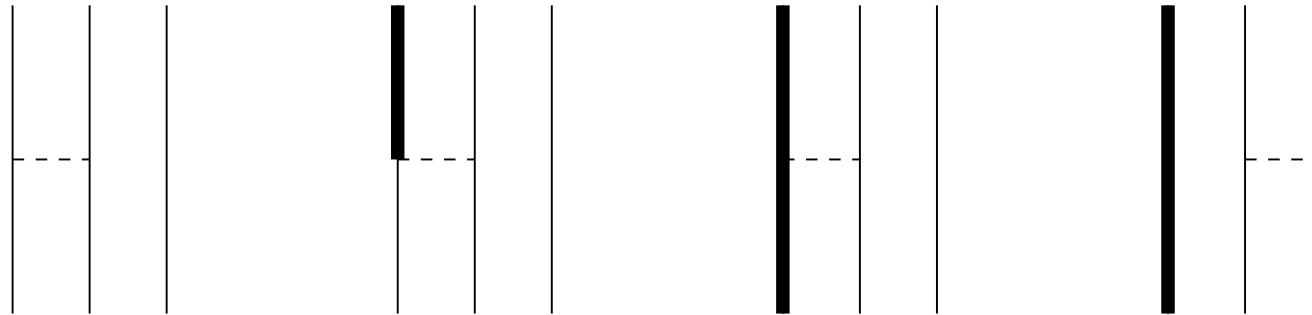


Core excitation (CX): extended Hilbert space

$$\mathcal{H} = \mathcal{H}_g \oplus \mathcal{H}_x$$



sector coupling by interaction



standard operator form of 3-body AGS equations

with $H_0 \rightarrow H_0 + h_A^{\text{int}}$

$$h_A^{\text{int}} |\mathcal{H}_a\rangle = (m_{A^*} - m_A) \delta_{ax} |\mathcal{H}_a\rangle$$

3-body AGS equations with core excitation

$$U_{\beta\alpha}^{ba} = \bar{\delta}_{\beta\alpha}\delta_{ba}G_0^{-1} + \sum_{\sigma} \sum_j \bar{\delta}_{\beta\sigma} T_{\sigma}^{bj} G_0 U_{\sigma\alpha}^{ja}$$

$$U_{0\alpha}^{ba} = \delta_{ba}G_0^{-1} + \sum_{\sigma} \sum_j T_{\sigma}^{bj} G_0 U_{\sigma\alpha}^{ja}$$

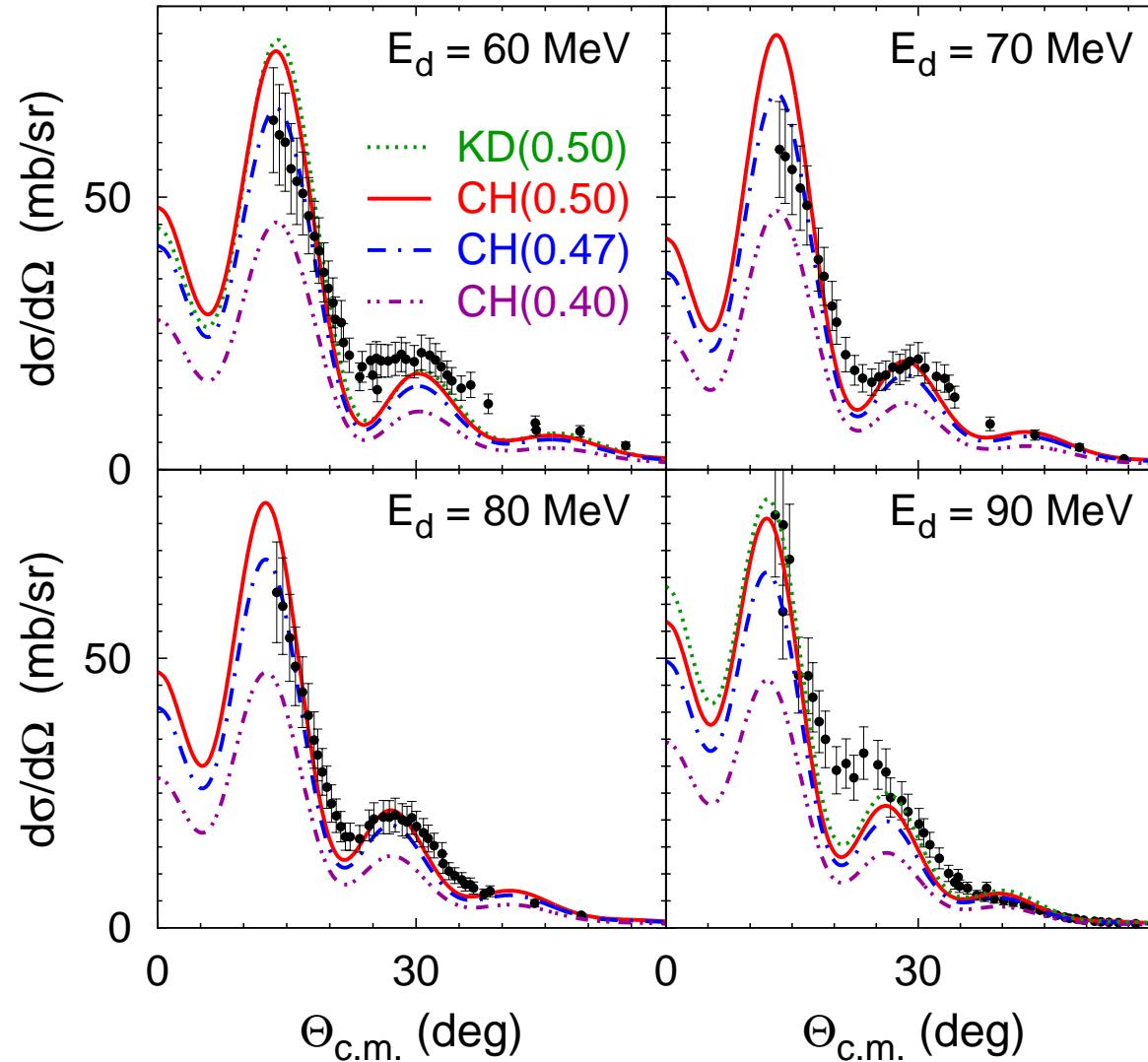
$$T_{\sigma}^{ba} = v_{\sigma}^{ba} + \sum_j v_{\sigma}^{bj} G_0 T_{\sigma}^{ja}$$

$$G_0 = (E + i0 - H_0)^{-1}$$

channel states $(E - H_0)|\phi_{\alpha}^a\rangle = \sum_j v_{\alpha}^{aj}|\phi_{\alpha}^j\rangle$

$$H_0|\mathbf{p}_{\alpha}\mathbf{q}_{\alpha}\rangle^a = [p_{\alpha}^2/2\mu_{\alpha} + q_{\alpha}^2/2M_{\alpha} + (\mathbf{m}_{A^*} - \mathbf{m}_A)\delta_{ax}]|\mathbf{p}_{\alpha}\mathbf{q}_{\alpha}\rangle^a$$

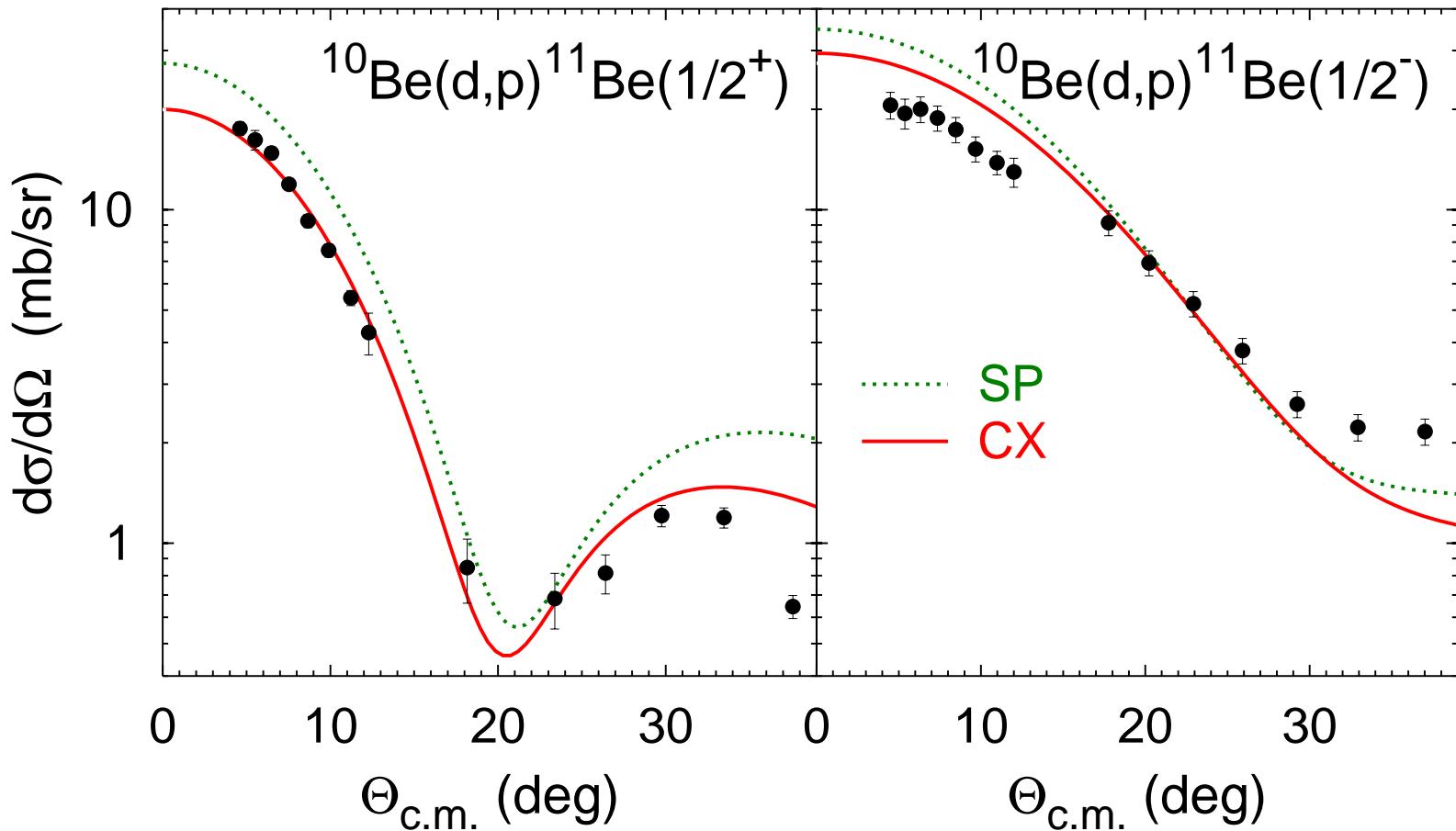
$^{24}\text{Mg}(\text{d},\text{d}')^{24}\text{Mg}(2^+)$ inelastic scattering



Rotational model for V_{NA} with $\beta_2 = 0.4...0.5$ [NPA 947, 173]

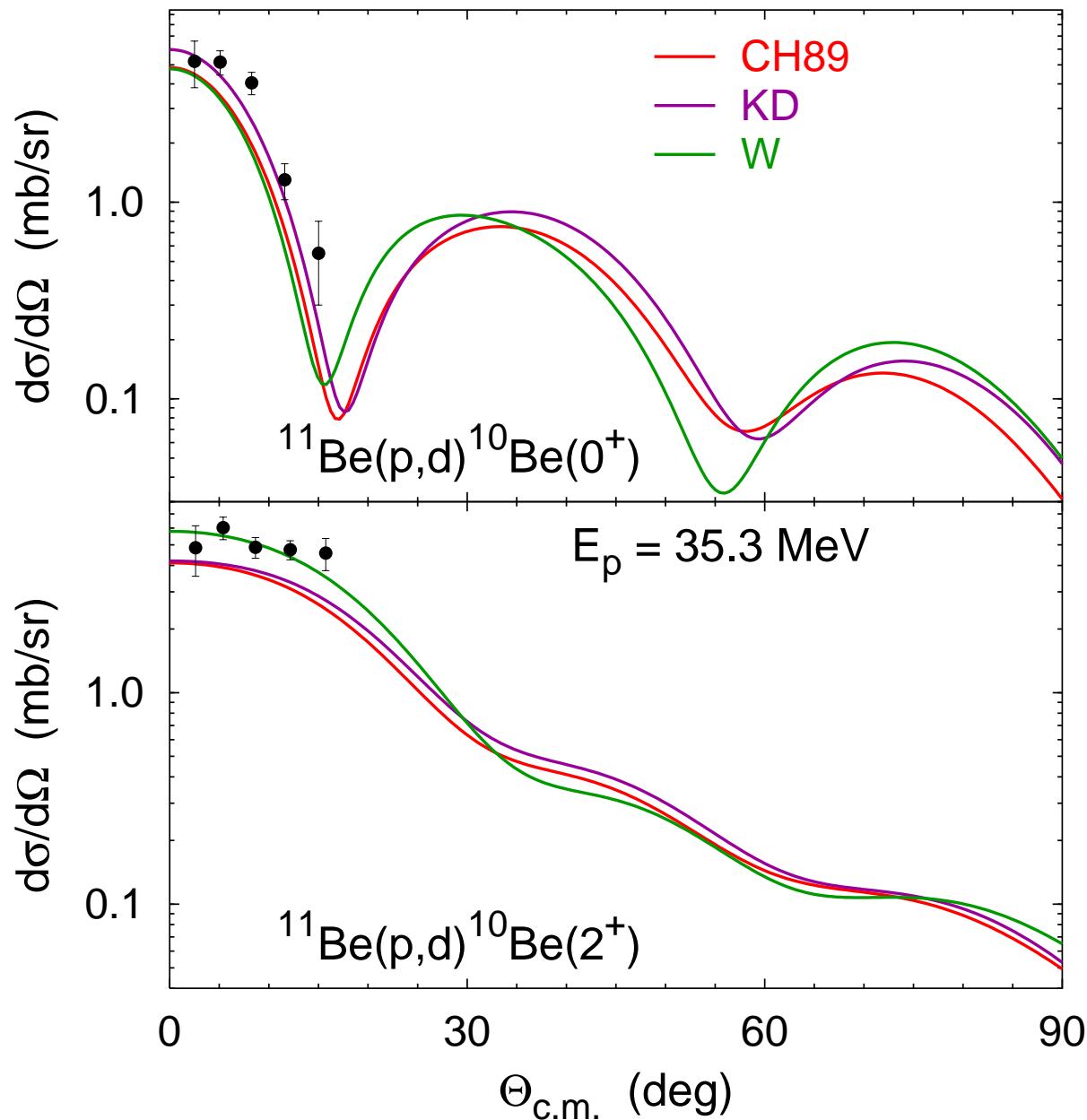
DWBA: $\beta_2 \sim 0.5 (p,p')$, $\beta_2 \sim 0.4 (d,d')$

CX effect in $^{10}\text{Be}(\text{d},\text{p})^{11}\text{Be}$ at 21.4 MeV

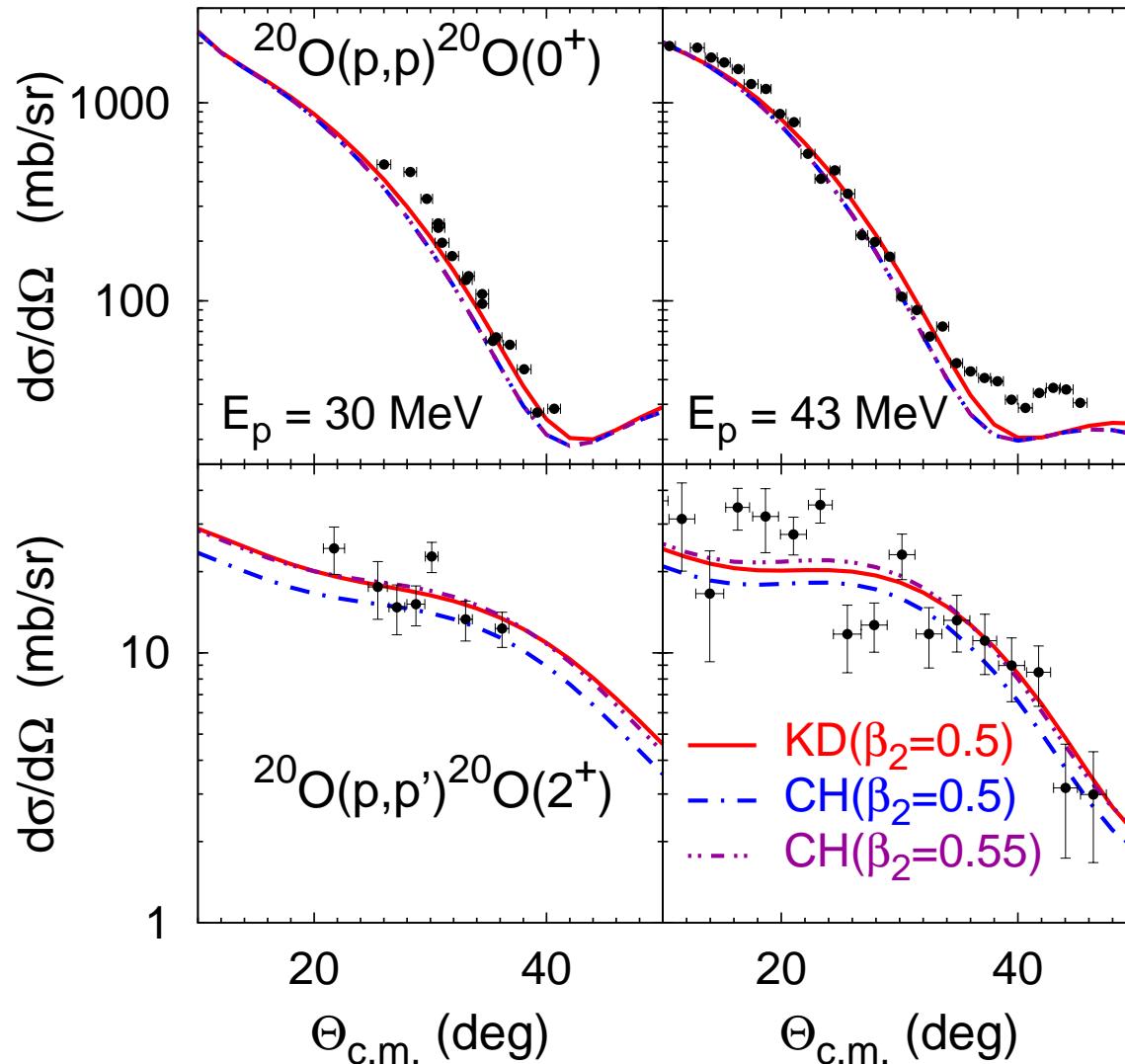


CH89, rotational model for V_{NA} with $\beta_2 = 0.67$
[PRC 91, 024607]

$^{11}\text{Be}(\text{p},\text{d})^{10}\text{Be}$: sensitivity to V_{NA}



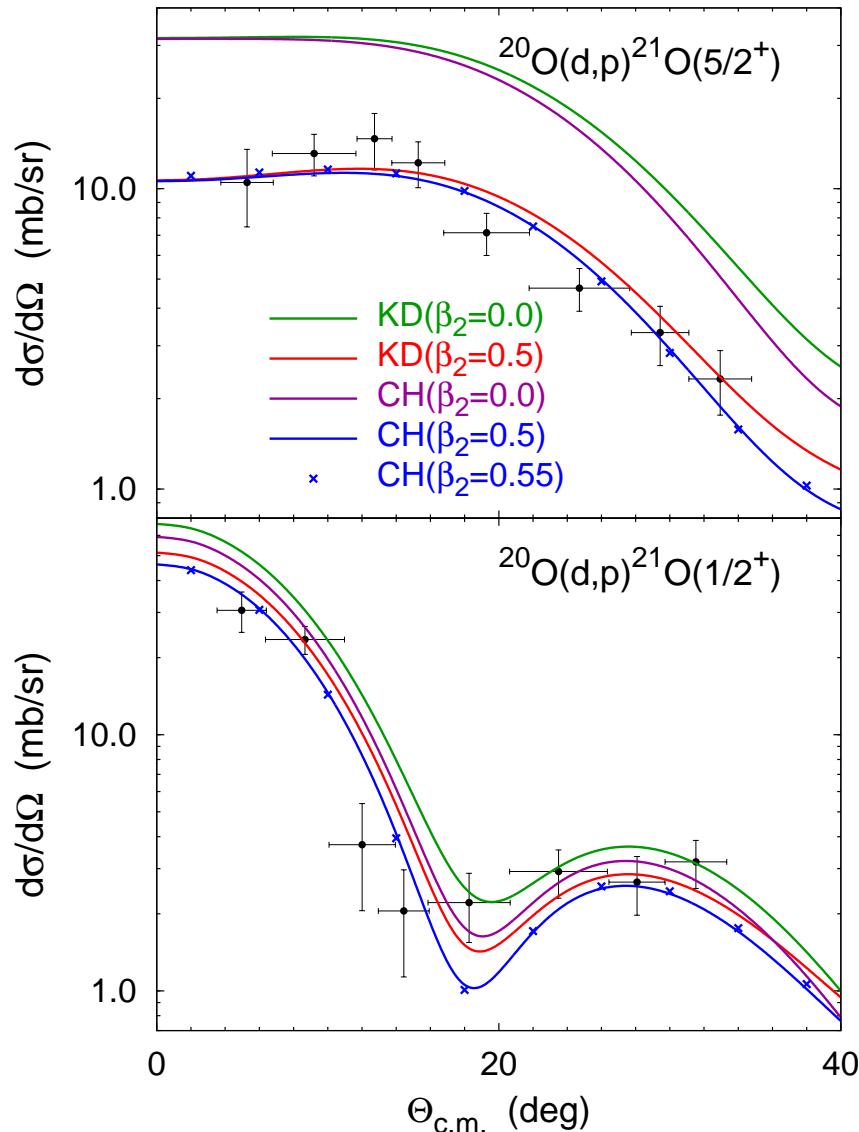
Potential test: N + ^{20}O



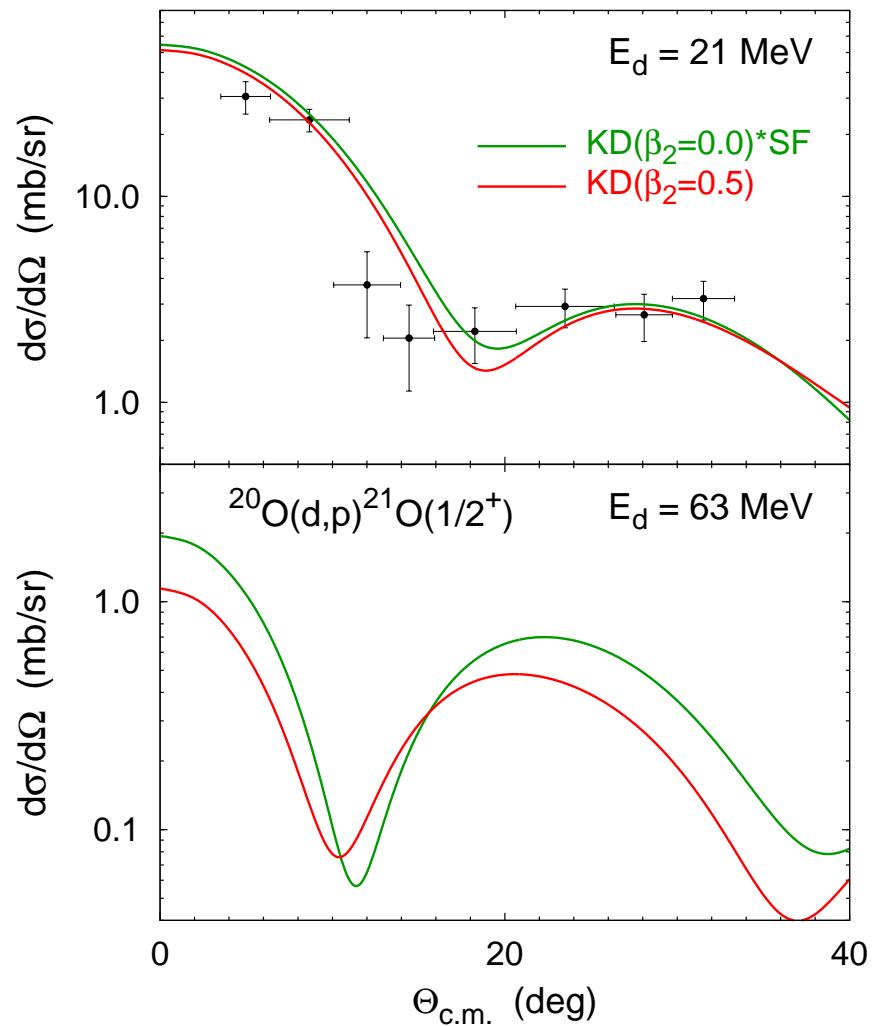
Vibrational model for V_{NA}

Shell-model SF for ^{21}O : $0.34(\frac{5}{2}^+), 0.82(\frac{1}{2}^+)$

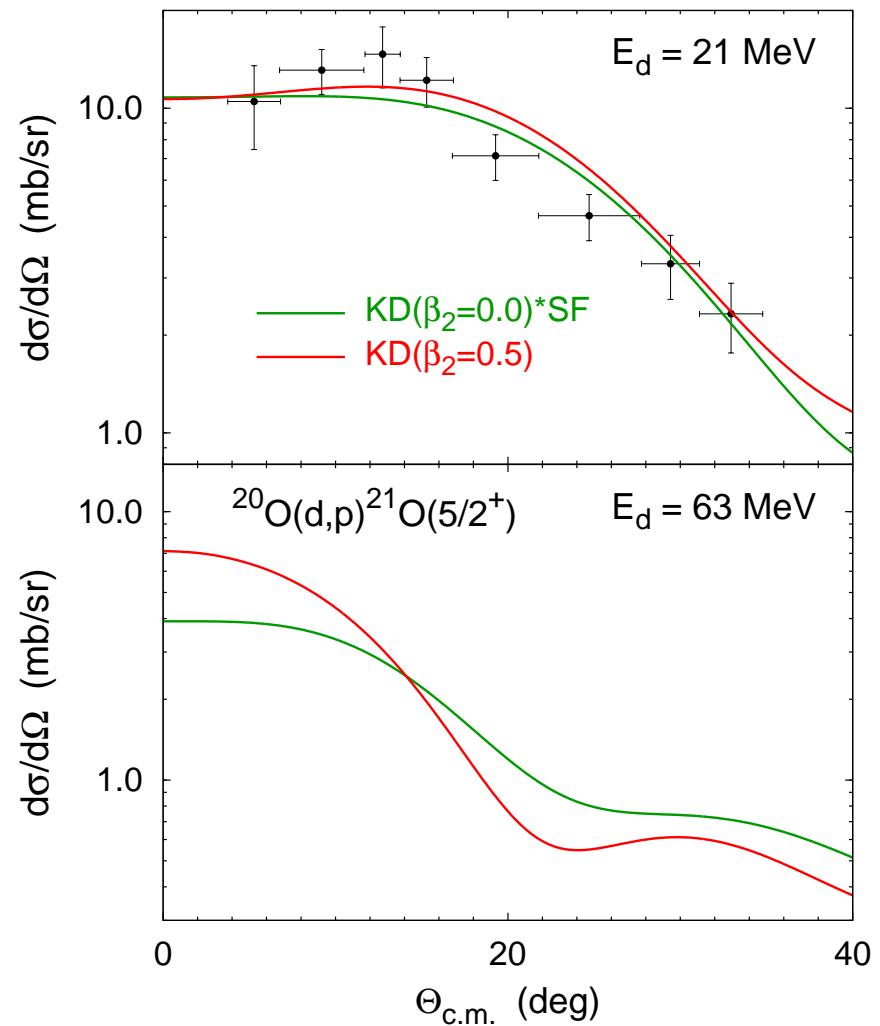
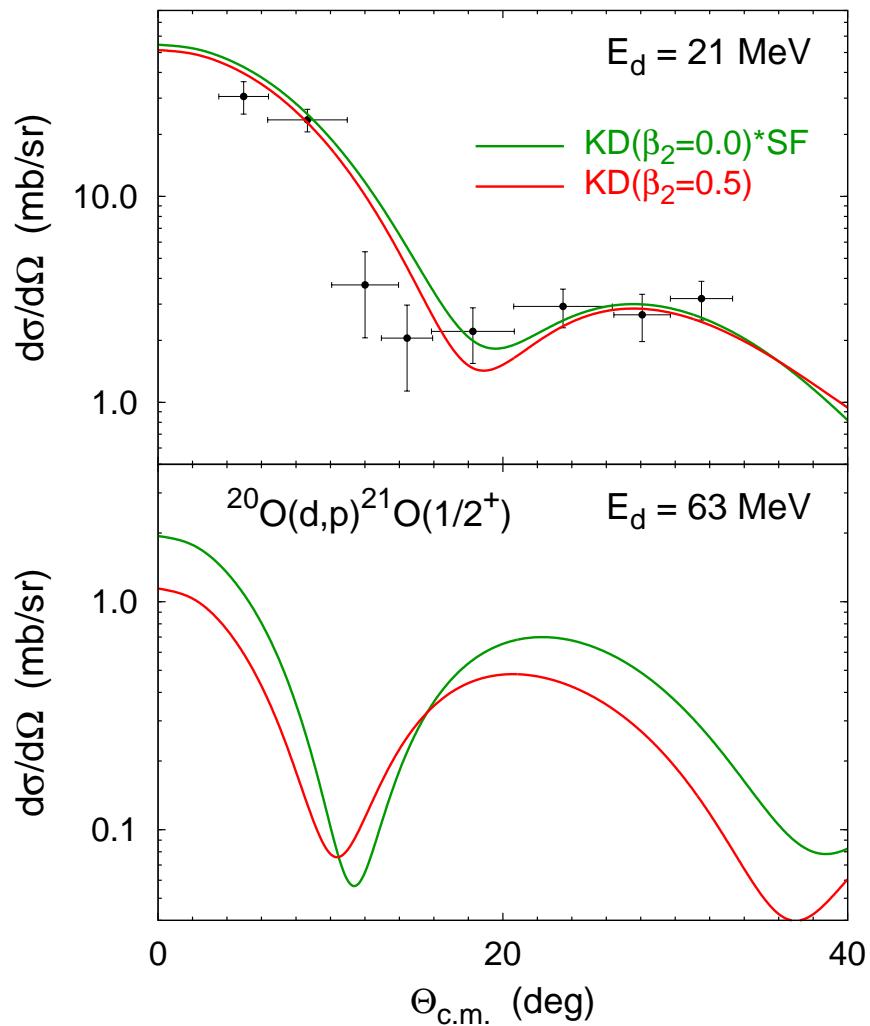
$^{20}\text{O}(\text{d},\text{p})^{21}\text{O}$ at 21 MeV



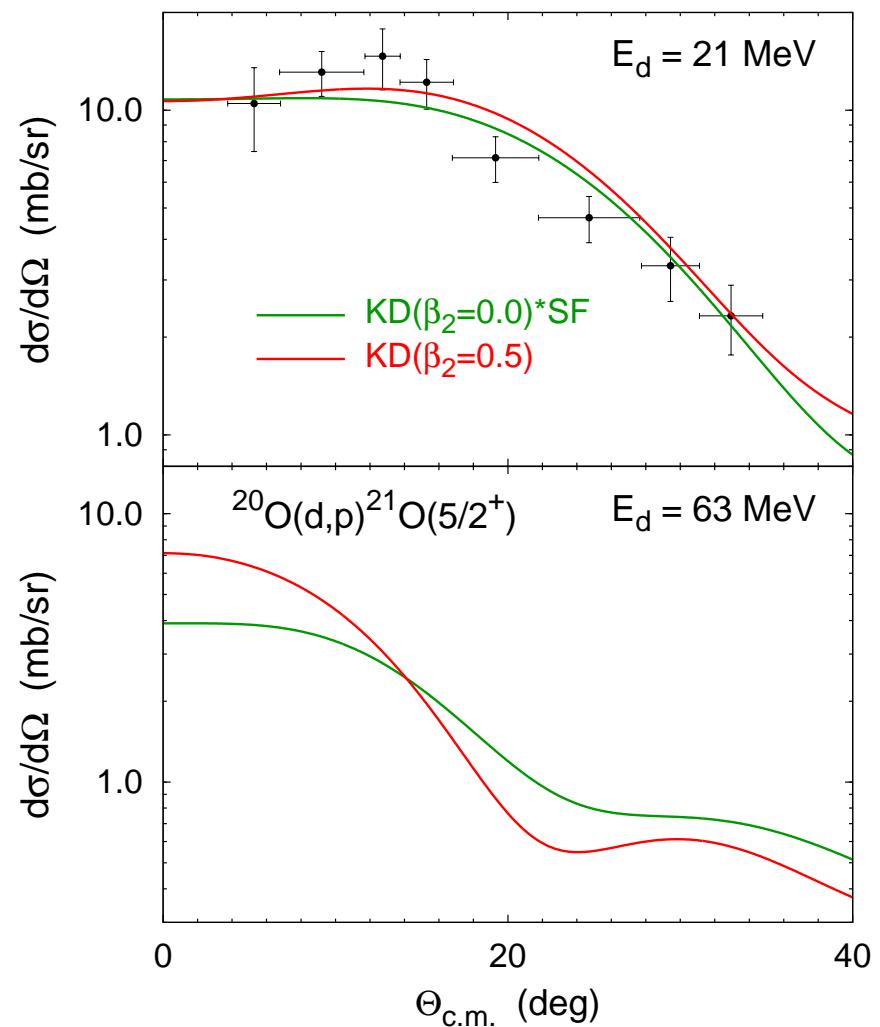
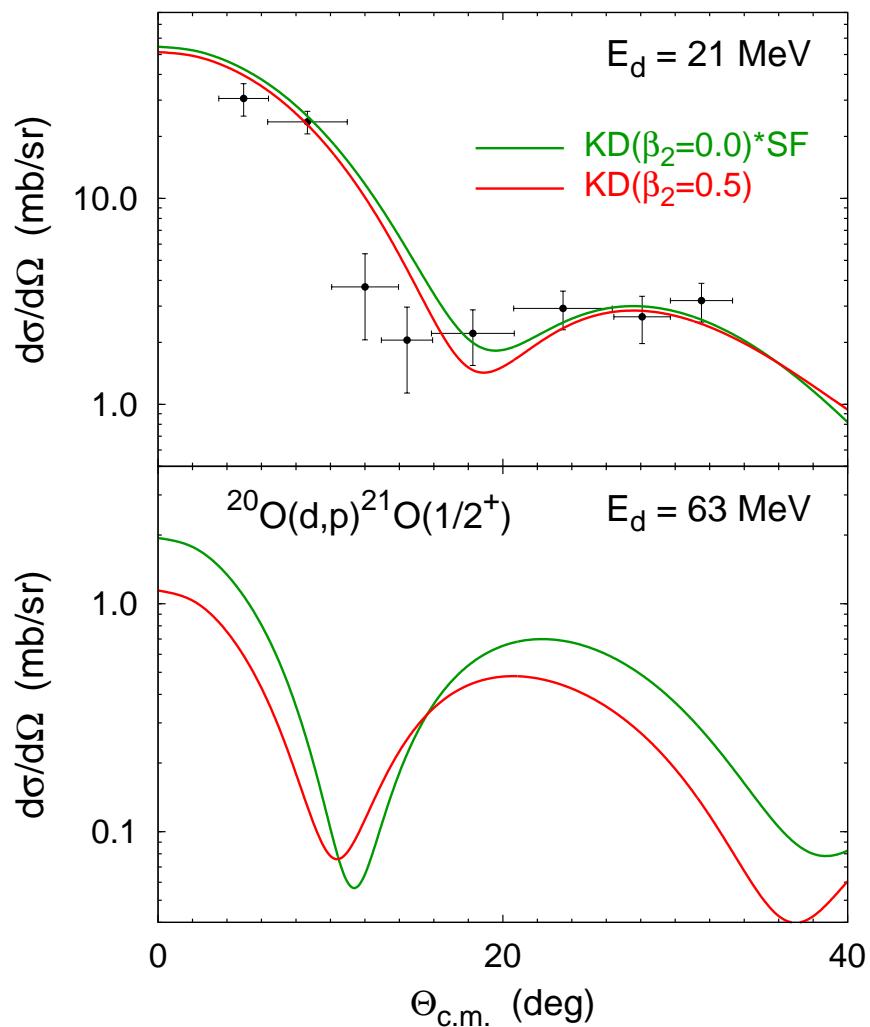
$^{20}\text{O}(\text{d},\text{p})^{21}\text{O}$: extracting SF?



$^{20}\text{O}(\text{d},\text{p})^{21}\text{O}$: extracting SF?



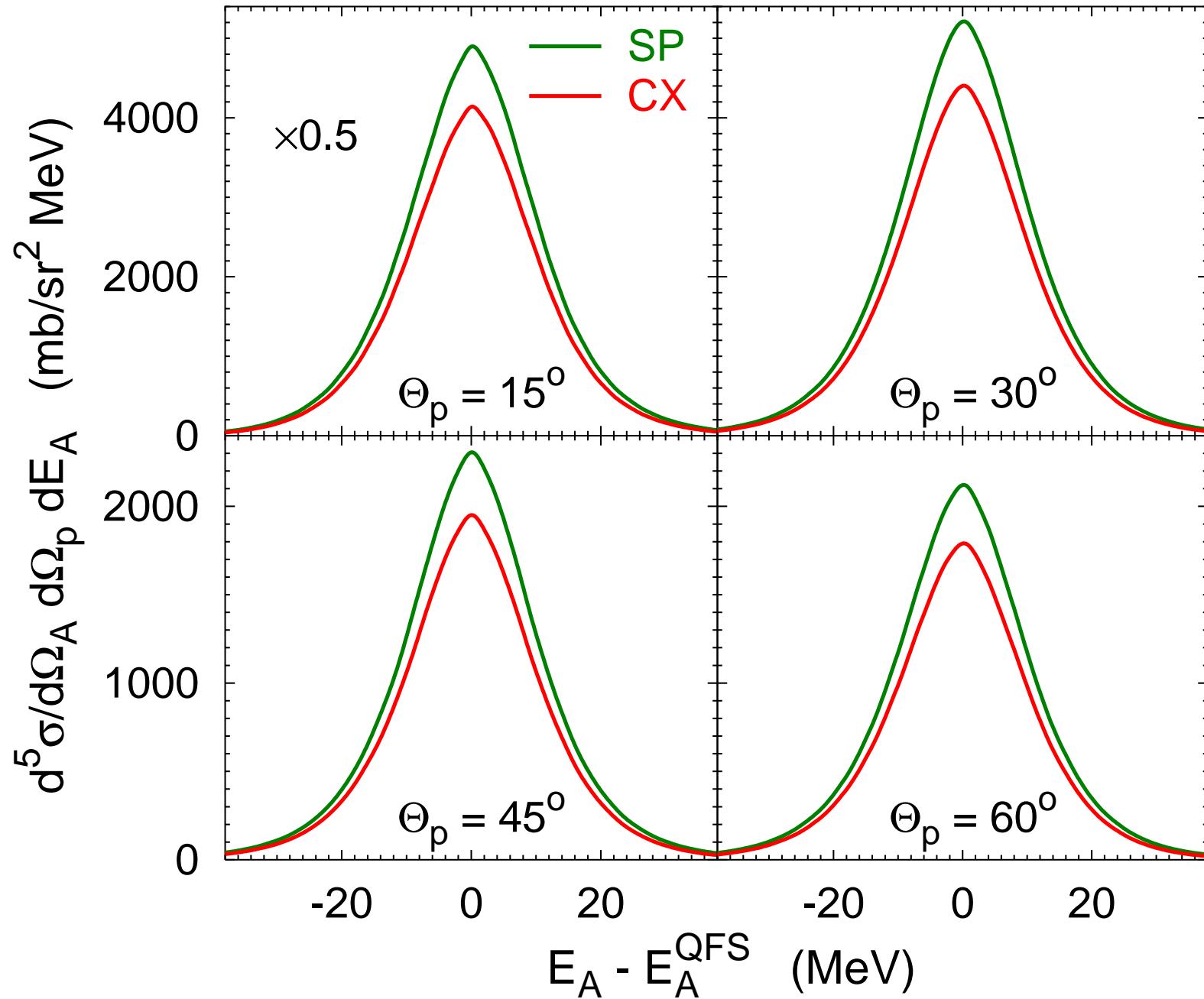
$^{20}\text{O}(\text{d},\text{p})^{21}\text{O}$: extracting SF?



SF = $\sigma_{\text{exp}}/\sigma_{\text{SP}}$ in general unreliable !

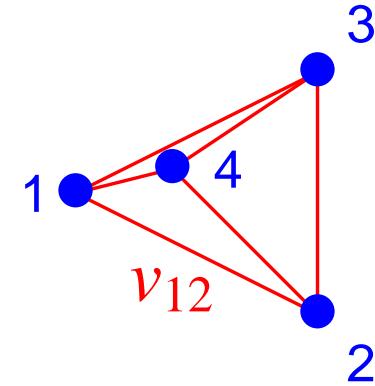
Faddeev/AGS: (V_{NA} - SF - data) compatibility check

$^{11}\text{Be}(\text{p},\text{pn})^{10}\text{Be}$ at 200 MeV/u near np QFS ($\Theta_A = 0^\circ$)



4N scattering

Hamiltonian $H_0 + \sum_{i>j} v_{ij}$



- Wave function:
Schrödinger equation (HH + Kohn VP, r -space)
[M. Viviani, A. Kievsky, L. E. Marcucci, S. Rosati, L. Girlanda]
- Wave function components:
Faddeev-Yakubovsky equations (r -space)
[R. Lazauskas, J. Carbonell]
- Transition operators:
Alt-Grassberger-Sandhas equations (p -space)
[AD, A. C. Fonseca]

4-body scattering: AGS equations

4-body transition operators

$$t_i = v_i + v_i G_0 t_i$$

$$G_0 = (E + i0 - H_0)^{-1}$$

$$U_{\gamma}^{jk} = G_0^{-1} \bar{\delta}_{jk} + \sum_i \bar{\delta}_{ji} t_i G_0 U_{\gamma}^{ik}$$

$$\mathcal{U}_{\beta\alpha}^{ji} = (G_0 t_i G_0)^{-1} \bar{\delta}_{\beta\alpha} \delta_{ji} + \sum_{\gamma k} \bar{\delta}_{\beta\gamma} U_{\gamma}^{jk} G_0 t_k G_0 \mathcal{U}_{\gamma\alpha}^{ki}$$

i, j, k : pairs (\equiv three-cluster (2+1+1) partitions)

α, β, γ : two-cluster (1+3 or 2+2) partitions

4-body scattering: AGS equations

4-body transition operators

$$t_i = v_i + v_i G_0 t_i$$

$$G_0 = (E + i0 - H_0)^{-1}$$

$$U_{\gamma}^{jk} = G_0^{-1} \bar{\delta}_{jk} + \sum_i \bar{\delta}_{ji} t_i G_0 U_{\gamma}^{ik}$$

$$\mathcal{U}_{\beta\alpha}^{ji} = (G_0 t_i G_0)^{-1} \bar{\delta}_{\beta\alpha} \delta_{ji} + \sum_{\gamma k} \bar{\delta}_{\beta\gamma} U_{\gamma}^{jk} G_0 t_k G_0 \mathcal{U}_{\gamma\alpha}^{ki}$$

i, j, k : pairs (\equiv three-cluster (2+1+1) partitions)

α, β, γ : two-cluster (1+3 or 2+2) partitions

wave function

$$|\Psi_{\alpha}\rangle = |\Phi_{\alpha}\rangle + \sum_{\gamma j k i} G_0 t_j G_0 U_{\gamma}^{jk} G_0 t_k G_0 \mathcal{U}_{\gamma\alpha}^{ki} |\phi_{\alpha}^i\rangle$$

$$|\Phi_{\alpha}\rangle = \sum_i |\phi_{\alpha}^i\rangle, \quad |\phi_{\alpha}^i\rangle = G_0 \sum_j \bar{\delta}_{ij} t_j |\phi_{\alpha}^j\rangle$$

4-body scattering amplitudes

two-cluster reactions:

$$\langle \Phi_\beta | T_{\beta\alpha} | \Phi_\alpha \rangle = \sum_{ji} \langle \phi_\beta^j | \mathcal{U}_{\beta\alpha}^{ji} | \phi_\alpha^i \rangle$$

three-cluster breakup:

$$\langle \Phi^j | T_\alpha^j | \Phi_\alpha \rangle = \sum_{\beta ki} \langle \Phi^j | \mathcal{U}_\beta^{jk} G_0 t_k G_0 \mathcal{U}_{\beta\alpha}^{ki} | \phi_\alpha^i \rangle$$

four-cluster breakup:

$$\langle \Phi_0 | T_{0\alpha} | \Phi_\alpha \rangle = \sum_{\beta jki} \langle \Phi_0 | t_j G_0 \mathcal{U}_\beta^{jk} G_0 t_k G_0 \mathcal{U}_{\beta\alpha}^{ki} | \phi_\alpha^i \rangle$$

[PRC 75, 014005; PRA 85, 012708]

Symmetrized AGS equations

$$t = v + v G_0 t$$

$$G_0 = (E + i\varepsilon - H_0)^{-1}$$

$$\textcolor{blue}{U}_j = P_j G_0^{-1} + P_j t G_0 \textcolor{blue}{U}_j$$

$$3+1 : \quad P_1 = P_{12} P_{23} + P_{13} P_{23}$$

$$2+2 : \quad P_2 = P_{13} P_{24}$$

$$\textcolor{red}{U}_{11} = (G_0 t G_0)^{-1} \zeta P_{34} + \zeta P_{34} \textcolor{blue}{U}_1 G_0 t G_0 \textcolor{red}{U}_{11} + \textcolor{blue}{U}_2 G_0 t G_0 \textcolor{red}{U}_{21}$$

$$\textcolor{red}{U}_{21} = (G_0 t G_0)^{-1} (1 + \zeta P_{34}) + (1 + \zeta P_{34}) \textcolor{blue}{U}_1 G_0 t G_0 \textcolor{red}{U}_{11}$$

$$\textcolor{red}{U}_{12} = (G_0 t G_0)^{-1} + \zeta P_{34} \textcolor{blue}{U}_1 G_0 t G_0 \textcolor{red}{U}_{12} + \textcolor{blue}{U}_2 G_0 t G_0 \textcolor{red}{U}_{22}$$

$$\textcolor{red}{U}_{22} = (1 + \zeta P_{34}) \textcolor{blue}{U}_1 G_0 t G_0 \textcolor{red}{U}_{12}$$

$$\zeta = -1 \text{ (+1) for fermions (bosons)}$$

basis states partially symmetrized

Scattering amplitudes: $E + i\varepsilon \rightarrow E + i0$

2-cluster reactions:

$$\textcolor{blue}{T}_{fi} = s_{fi} \langle \phi_f | \textcolor{red}{U}_{fi} | \phi_i \rangle$$

$$|\phi_j\rangle = G_0 t P_j |\phi_j\rangle$$

$$|\Phi_j\rangle = (1 + P_j) |\phi_j\rangle$$

3-cluster breakup/recombination:

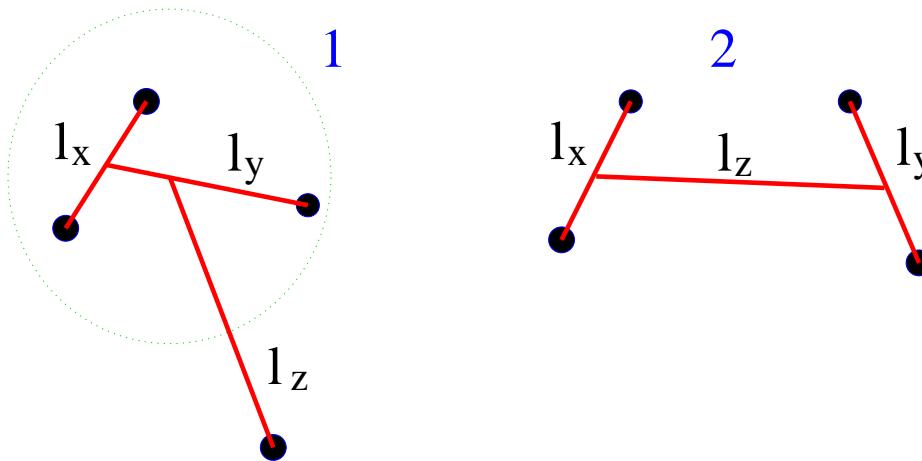
$$T_{3i} = s_{3i} \langle \phi_3 | [(1 + \zeta P_{34}) \textcolor{green}{U}_1 G_0 t G_0 \textcolor{red}{U}_{1i} + \textcolor{green}{U}_2 G_0 t G_0 \textcolor{red}{U}_{2i}] | \phi_i \rangle$$

4-cluster breakup/recombination:

$$\begin{aligned} \textcolor{blue}{T}_{4i} = & s_{4i} \{ \langle \phi_4 | [1 + (1 + P_1) \zeta P_{34}] (1 + P_1) t G_0 \textcolor{green}{U}_1 G_0 t G_0 \textcolor{red}{U}_{1i} | \phi_i \rangle \\ & + \langle \phi_4 | (1 + P_1) (1 + P_2) t G_0 \textcolor{green}{U}_2 G_0 t G_0 \textcolor{red}{U}_{2i} | \phi_i \rangle \} \end{aligned}$$

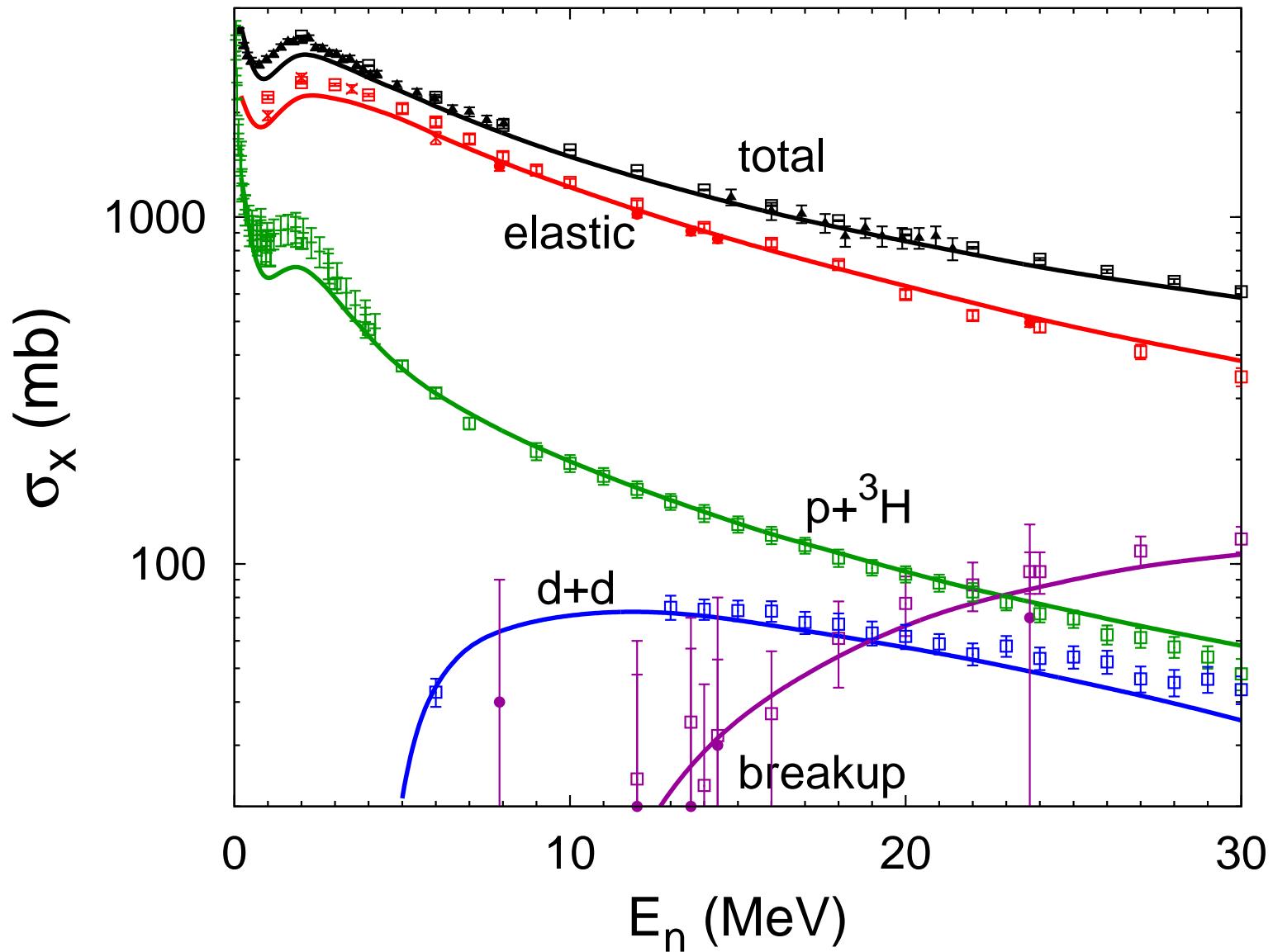
Solution of 4N AGS equations

$$U_{12}|\phi_2\rangle = G_0^{-1}P_2|\phi_2\rangle - P_{34}U_1G_0tG_0U_{12}|\phi_2\rangle + U_2G_0tG_0U_{22}|\phi_2\rangle$$



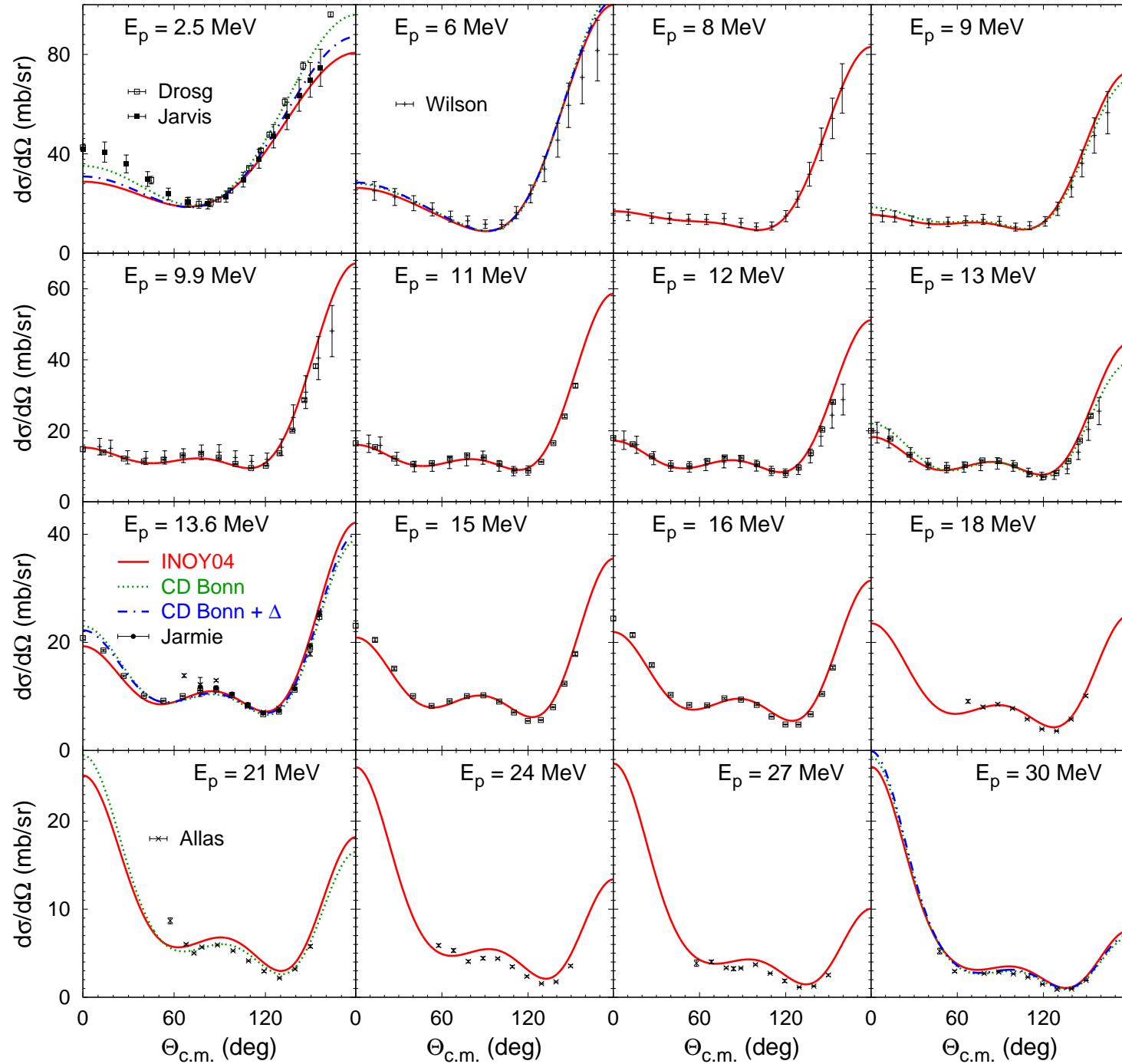
- momentum-space partial-wave basis
 $|k_x k_y k_z [l_z (\{l_y [(l_x S_x) j_x s_y] S_y\} J_y s_z) S_z] JM, [(T_x t_y) T_y t_z] TM_T \rangle_1$
 $|k_x k_y k_z [l_z \{(l_x S_x) j_x [l_y (s_y s_z) S_y] j_y\} S_z] JM, [T_x (t_y t_z) T_z] TM_T \rangle_2$
- large system (up to 30000) of coupled 3-variable integral equations with integrable singularities
- Coulomb interaction: screening and renormalization
[PRC 75, 014005, PRL 98, 162502]

$n + {}^3\text{He}$ total and partial cross sections

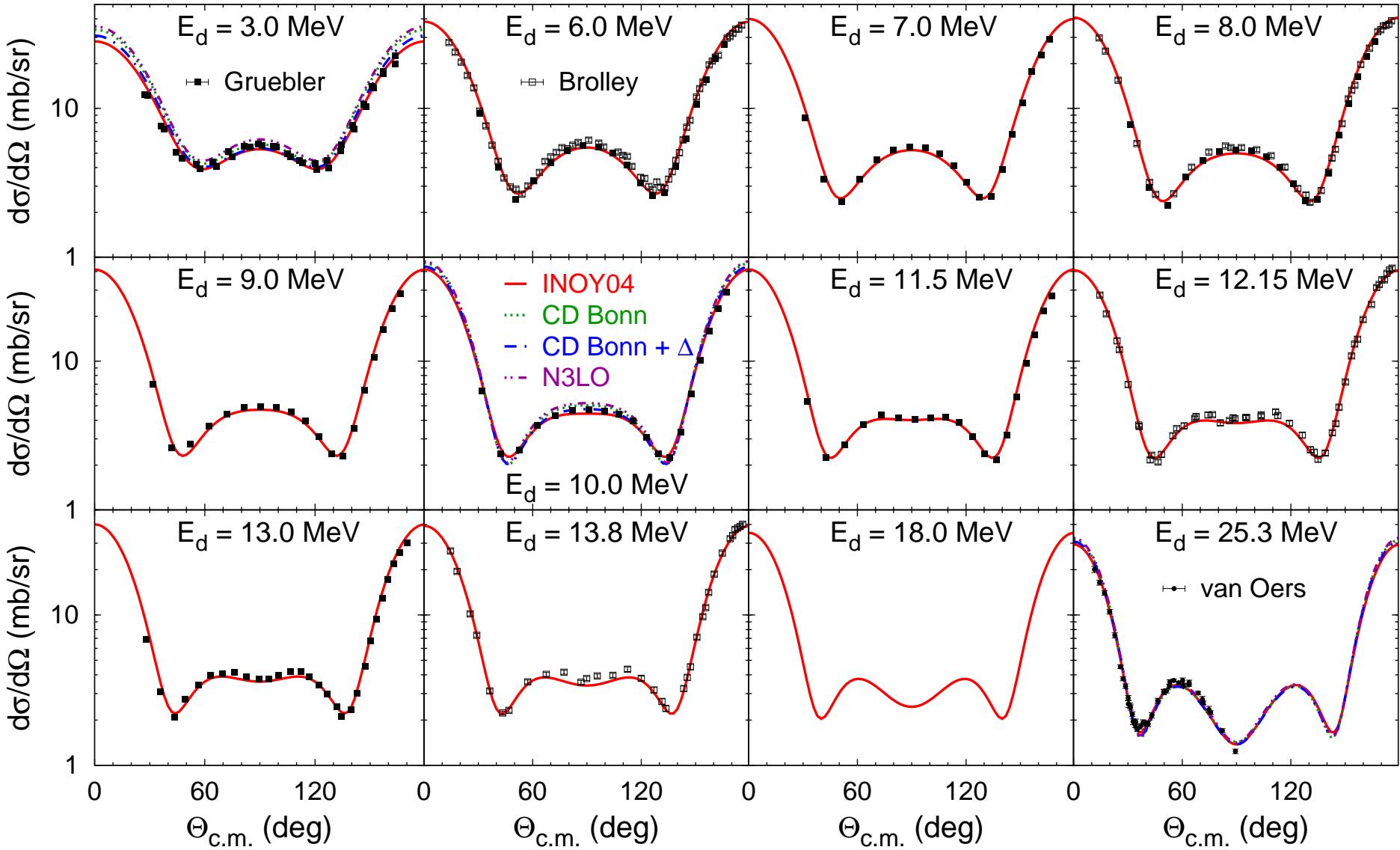


[PRL 113, 102502; PRC 90, 044002]

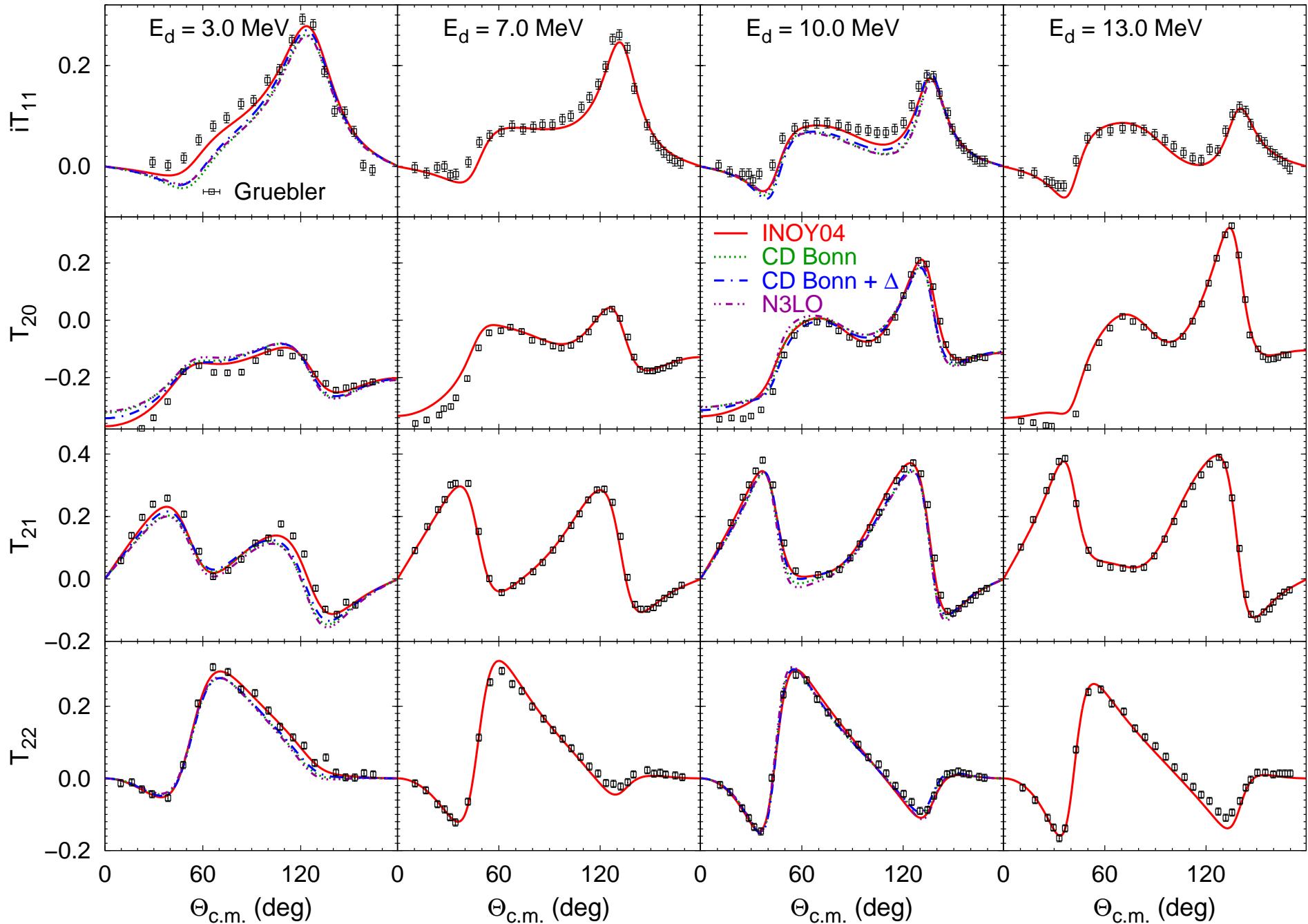
Charge exchange reaction ${}^3\text{H}(p,n){}^3\text{He}$



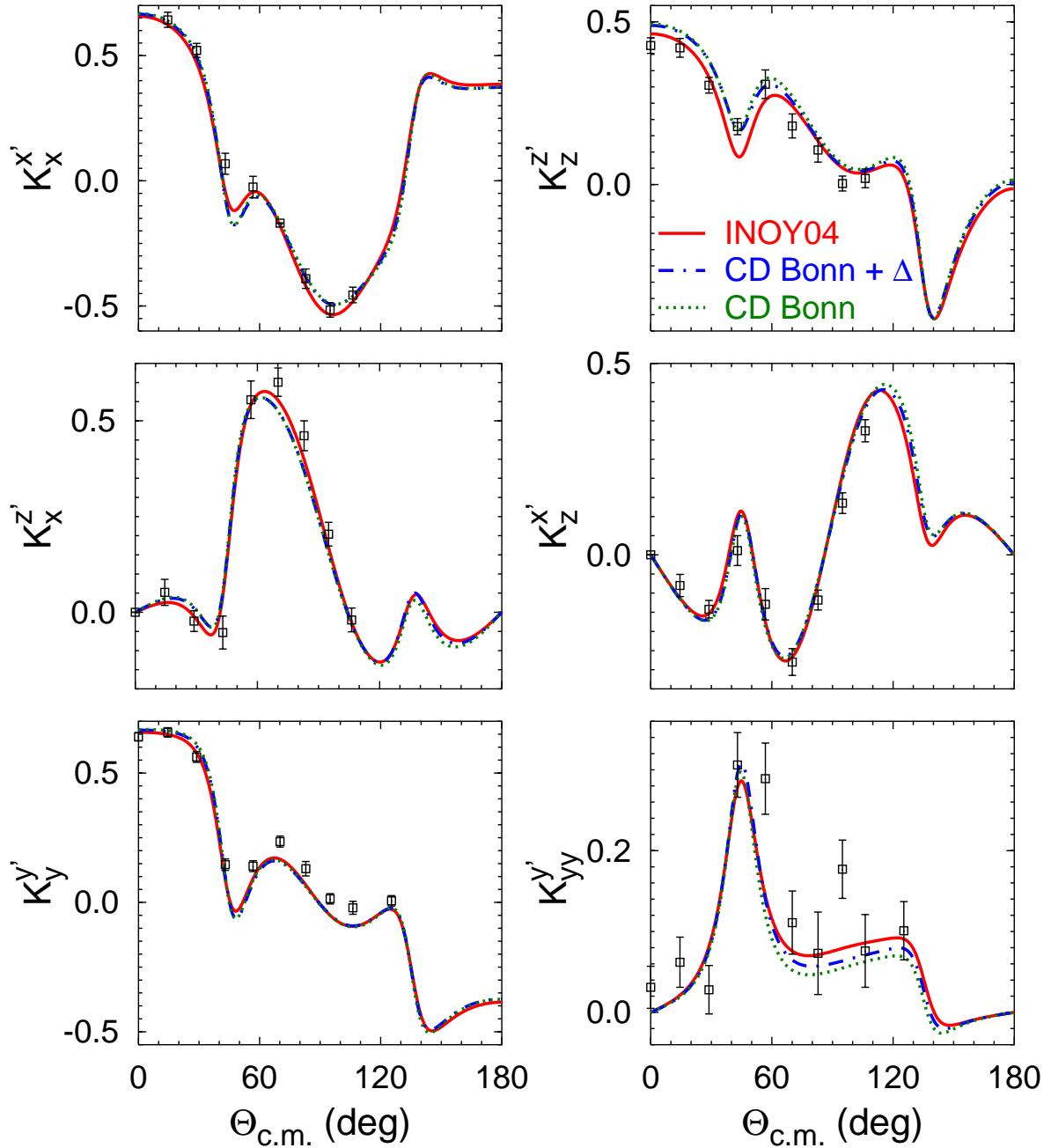
Transfer reaction $^2\text{H}(d, p)^3\text{H}$



Transfer reaction $^2\text{H}(\vec{d}, p)^3\text{H}$: analyzing powers



Spin transfer in $^2\text{H}(\vec{d}, \vec{n})^3\text{He}$ at 10 MeV



Few-body nuclear reactions

- 3-body AGS equations:
extension including core excitation
- complicated CX effects in transfer reactions,
no simple relation to SF
- 4-body AGS equations
— (p,n), (d,p) and (d,n) reactions in 4N system