

Resonance phenomena: from compound nucleus decay to proton  
radioactivity  
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- 1 Introduction
- 2 Resonance properties
  - Single model
  - Gamow complex energy
  - $P_{well}$  and phase shift
- 3 Symmetry dependence of mean field potential
- 4 Summary

# Introduction

## Intro

- The low-lying excited states of exotic nuclei near the drip line are in the continuum.
- Even the ground states are resonances when beyond the drip line.
- Resonances play an important role in the structure of exotic nuclei.

# Resonance properties

## single model

### One-dimensional model

Consider a simple one-dimensional potential with some important features:

- A potential well with a finite depth  $V_0$ . ( $x < x_0$ )
- A barrier with a height of  $V_b$ . ( $x_0 < x < x_b$ )
- A continuum  $\rightarrow$  the potential **out of the barrier** equal to zero. ( $x > x_b$ )
- The solutions in different areas respectively:

$$\psi(x) = \mathcal{A}\sin(k_0x) + \mathcal{B}\cos(k_0x) = \mathcal{A}\sin(k_0x) \quad x < x_0 \quad (1)$$

$$\psi(x) = \mathcal{E}\exp(\kappa_\infty x) + \mathcal{F}\exp(-\kappa_\infty x) = \mathcal{F}\exp(-\kappa_\infty x) \quad x > x_b \quad (2)$$

$$\psi(x) = \mathcal{C}\exp(\kappa_b x) + \mathcal{D}\exp(-\kappa_b x) \quad x_0 < x < x_b \quad (3)$$

with  $k_0 = \sqrt{2m(E + V_0)}/\hbar$ ,  $\kappa_\infty = \sqrt{2m|E|}/\hbar$ ,  $\kappa_b = \sqrt{2m(V_b - E)}/\hbar$ .

# Resonance properties

single model

## Solution of the *Schrödinger* equation

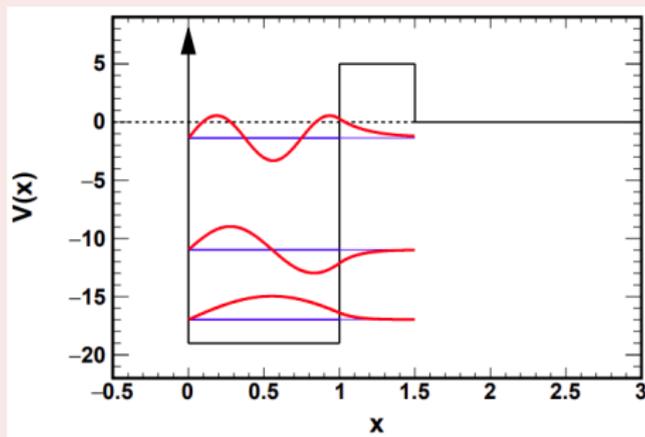


Figure 1: The solution with the example simple potential mentioned before

- blue horizontal lines: bound energies
- red lines: the corresponding wavefunctions

# Resonance properties

from bound to resonance

What if we decrease the depth of the well  $V_0$ ?  
→ The least-bound level will cross zero energy, what kind of state?

# Resonance properties

from bound to resonance

## Quantum tunneling

- $0 < E < V_0$
- Each time the nucleon attacks the barrier, there is a probability of penetrating the barrier.
- The probability for the particle staying behind the barrier is reduced by a constant decreasing exponentially with time.

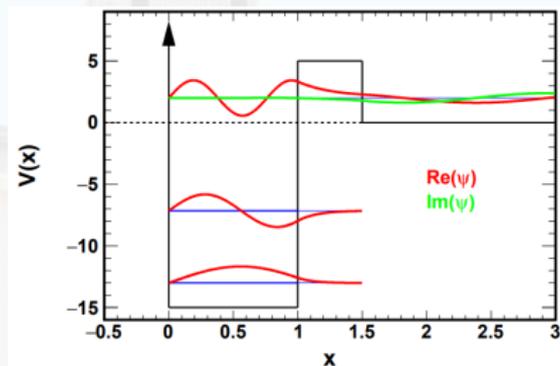


Figure 2: With a shallower well, the least-bound state becomes a resonance state.

# Resonance properties

## Gamow complex energy

### Exponential-decaying solutions of *Schrödinger* equation (by Gamow)

- time-dependent *Schrödinger* equation

$$i\hbar \frac{\partial}{\partial t} \Psi(x, t) = \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(r) \right] \Psi(x, t)$$

- use a complex energy  $E = E_r - i\Gamma/2$ , the modulus squared of the wave function.

$$|\Psi(x, t)|^2 = \exp\left(\frac{-\Gamma t}{\hbar}\right) |\psi(x)|^2$$

- by Heisenberg's uncertainty principle we can define the width and lifetime.

$$\Delta E_r \Delta t \sim \hbar$$

- the state has a finite lifetime  $\tau = \Delta t = \frac{\hbar}{\Gamma}$

## Resonance properties

$P_{well}$  and phase shift

### The probability of being in the well

- Considering the potential in Figure.2, we define the probability as

$$P_{well} = \int_0^{x_0} |\psi(x)|^2 dx$$

- In addition, the wavefunction in the exterior region( $x > x_b$ )

$$\psi(x) = \frac{\sqrt{\mathcal{A}^2 + \mathcal{B}^2}}{2i} \exp(-\delta) [\exp(ik_\infty x + 2\delta) - \exp(-ik_\infty x)]$$

- This leads to a phase shift of  $2\delta$  between the incoming and outgoing waves.

## Resonance properties

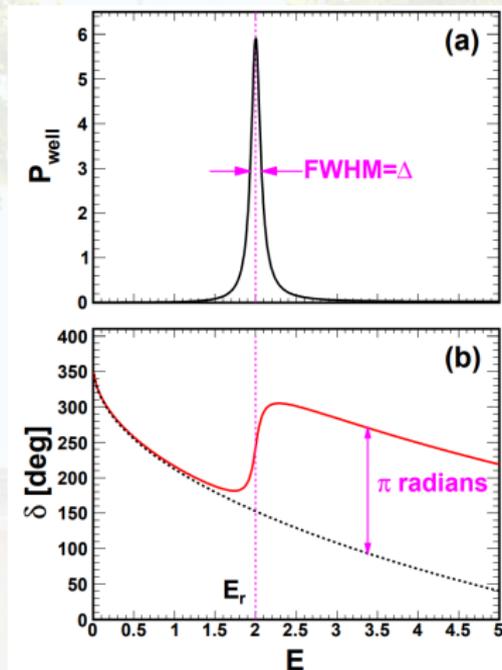
 $P_{well}$  and phase shift

Figure 3: Energy dependence of the phase shift and  $P_{well}$  (the relative probability). At the resonance energy, the phase shift shows a sharp jump of  $\pi$  radians.

# Resonance properties

$P_{well}$  and phase shift

## Breit-Wigner resonance

- The relative probability of getting a particle inside of the potential well has a strong **peak**  $\rightarrow$  **resonance**.
- For narrow peaks, the centroid of a peak has the same value as  $E_r$
- The peak's shape is given by Breit-Wigner form:

$$P_{well} \propto \frac{\Gamma^2}{(E - E_r)^2 + (\Gamma/2)^2}$$

- $\Gamma$ : the **full width at half maximum (FWHM)** of the peak.
- The ' $\Gamma$ ' here is the same as that of the imaginary part of complex energy ( $-\Gamma/2$ ).
- The phase shift shows a sudden change of magnitude  $\pi$  radians relative to the slowly varying background.

## Symmetry dependence of mean field potential

We take the different kinds of nucleons into consideration.

the modification of the potential deriving from the two nucleon types

- keep the same number of nucleons but change the ratios of p to n
- the depth of the mean field

$$V_0 = V' + v_{sym} \frac{N - Z}{A} \text{protons} \quad (4)$$

$$= V' - v_{sym} \frac{N - Z}{A} \text{neutrons} \quad (5)$$

- This kind of symmetry dependence contributes about 50% of the symmetry energy.(emperical)
- The binding energy

$$E_{binding}(Z, A) = a_{vol}A - a_{sur}A^{2/3} - a_{Coul} \frac{Z^2}{A^{1/3}} - a_{asy} \frac{(N - Z)^2}{A}$$

involves the volume, surface, Coulomb, and symmetry terms.

## Symmetry dependence of mean field potential

the importance of the symmetry force for the single-particle energies

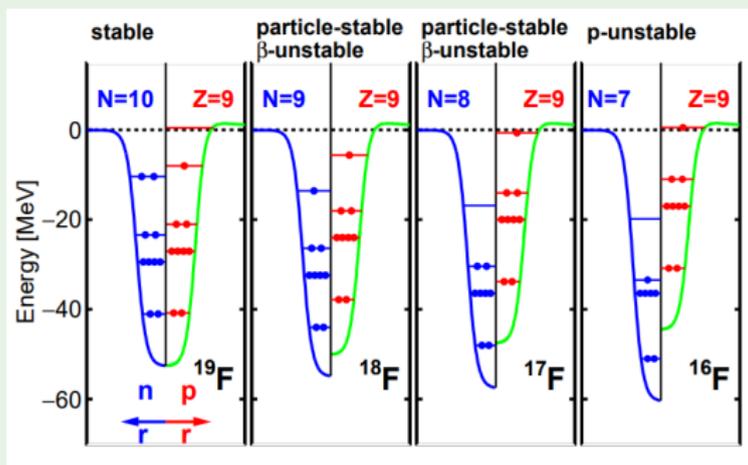


Figure 4: Schematic showing the evolution of neutron and proton single-particle levels with mass number  $A$  for fluorine isotopes.

- With the decrease of the number of neutrons, the energy level of protons increase.  $\rightarrow$  making a proton resonance.

# Summary

## Main theoretical part of this article

- 1 A resonance state will decay with an exponential time distribution with a lifetime of  $\hbar/\Gamma$ .
  - 2 Resonance can be detected in scattering experiments in that the cross section corresponds to the  $P_{well}$  with a FWHM of  $\Gamma$ .
  - 3 A resonance has a sharp variation (which is a jump of  $\pi$  radians) on the phase shift when the energy scans across the resonance energy.
  - 4 We can modified the number of one type of nucleons thus changing the mean field of each nucleon, making a proton or a neutron resonance.
- Next time we will cover the left part of this article.

# Thank You!