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Dynamical eikonal

对于projectile，考虑其两体结构由core和fragment构成。

弹核内部的自由度由内部哈密顿量描述：

$$H_0 = \frac{p^2}{2\mu_{cf}} + V_{cf}(\mathbf{r}), \quad H_0\phi_0 = E_0\phi_0$$

三体哈密顿量的形式：

$$H = \frac{P^2}{2\mu} + H_0 + V_{cT} \left(\mathbf{R} - \frac{m_f}{m_P} \mathbf{r} \right) + V_{fT} \left(\mathbf{R} + \frac{m_c}{m_P} \mathbf{r} \right)$$

目的是求解三体的薛定谔方程： $H\Psi = E_T\Psi$

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Eikonal ansatz:

$$\Psi(\mathbf{R}, r) = e^{iKZ} \hat{\Psi}(\mathbf{R}, r)$$

K为相对波矢:

$$\frac{\hbar^2 K^2}{2\mu} = E_T - E_0$$

代入三体哈密顿量之后得到约化后的薛定谔方程:

$$\frac{\hbar^2}{2\mu} \Delta_R \hat{\Psi} + i\hbar v \frac{\partial \hat{\Psi}}{\partial Z} = \left(H_0 + V_{cT} + V_{fT} - E_0 \right) \hat{\Psi}$$

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$$\frac{\hbar^2}{2\mu} \Delta_R \hat{\Psi} + i\hbar v \frac{\partial \hat{\Psi}}{\partial Z} = \left(H_0 + V_{cT} + V_{fT} - E_0 \right) \hat{\Psi}$$

传统的eikonal近似包含两项：1、忽略二阶导数。2、绝热近似。

$$i\hbar v \frac{\partial \hat{\Psi}}{\partial Z} = (V_{cT} + V_{fT}) \hat{\Psi} \quad \hat{\Psi} = \exp \left[\int_{-\infty}^z \frac{U_{cT} + U_{fT}}{i\hbar v} dz \right]$$

本质上还是两体问题

直接完全忽略了projectile的自由度

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$$\frac{\hbar^2}{2\mu} \Delta_R \hat{\Psi} + i\hbar v \frac{\partial \hat{\Psi}}{\partial Z} = \left(H_0 + V_{cT} + V_{fT} - E_0 \right) \hat{\Psi}$$

Dynamical eikonal只忽略二阶导数

$$i\hbar v \frac{\partial \hat{\Psi}}{\partial Z} = \left(H_0 + V_{cT} + V_{fT} - E_0 \right) \hat{\Psi}$$

引入类似于时间的变量: $v dt = dz$

“... is formally equivalent to the semiclassical time-dependent Schrodinger equation with straight-line trajectories [1].”

$$i\hbar \frac{\partial \hat{\Psi}}{\partial t} = \left(H_0 + V_{cT} + V_{fT} - E_0 \right) \hat{\Psi}$$

[1] Baye, D., Capel, P., & Goldstein, G. (2005). Collisions of halo nuclei within a dynamical eikonal approximation. *Physical review letters*, 95(8), 082502.

Dynamical eikonal

直接计算 T 矩阵:

$$T_{fi} = \left\langle e^{i\mathbf{K}\cdot\mathbf{R}} \phi_0(\mathbf{r}) \left| V_{cT} + V_{fT} \right| \Psi(\mathbf{R}, \mathbf{r}) \right\rangle$$

$$T'_{fi} = \left\langle e^{i\mathbf{K}\cdot\mathbf{R}} \chi^{(-)}(\mathbf{r}) \left| V_{cT} + V_{fT} \right| \Psi(\mathbf{R}, \mathbf{r}) \right\rangle$$

转化为做二维傅立叶变换:

$$\begin{aligned} T_{fi} &= i\hbar v \left\langle e^{i\mathbf{K}\cdot\mathbf{R}} \phi_0(\mathbf{r}) \left| e^{i\mathbf{K}Z} \frac{\partial}{\partial Z} \hat{\Psi}(\mathbf{R}, \mathbf{r}) \right. \right\rangle \\ &\approx i\hbar v \int d\mathbf{R} e^{-i\mathbf{q}\cdot\mathbf{b}} \frac{\partial}{\partial Z} \left\langle \phi_0(\mathbf{r}) \left| \hat{\Psi}(\mathbf{R}, \mathbf{r}) \right. \right\rangle \end{aligned}$$

Hussein-McVoy 模型

NEB的形式为:

$$\left. \frac{d^2\sigma}{dE_b d\Omega_b} \right|_{\text{NEB}}^{\text{HM}} = -\frac{2}{\hbar v_a} \rho_b(E_b) \left\langle \varphi_x^{\text{HM}} \left| W_x \right| \varphi_x^{\text{HM}} \right\rangle$$

x-A波函数的形式:

$$\varphi_x^{\text{HM}}(\mathbf{r}_x) = \langle \mathbf{r}_x | \varphi_x^{\text{HM}} \rangle = \left\langle \mathbf{r}_x \chi_b^{(-)} \left| \chi_a^{(+)} \phi_a \right. \right\rangle$$

不需要求解耦合方程/计算source term

对整个模型空间进行了简化

Hussein-McVoy 模型

$$\sigma_{\text{NEB}}^{\text{EHM}} = \frac{2}{v_a} (2\pi)^3 \frac{E_x}{\hbar k_x} \int d^3 \mathbf{r}_b d^3 \mathbf{r}_x \left| \phi_a(\mathbf{r}_{bx}) \right|^2 \\ \times \left| S_{bA}(b_b) \right|^2 \left[1 - \left| S_{xA}(b_x) \right|^2 \right].$$

NEB的形式为：

“This equation has an appealing and intuitive form: the integrand contains the product of the probabilities for the core being elastically scattered by the target, $|S_{bA}|^2$, times the probability of the valence particle being absorbed, $(1 - |S_{xA}|^2)$.” [2]

“depends only on the asymptotic properties, this is, the S matrices, of the interaction of b and x with the target. There is no sensitivity to the wave- functions in the interaction region. This is a result of the eikonal approximation...”

[2] Gómez-Ramos, M., Gómez-Camacho, J., Lei, J. *et al.* The Hussein–McVoy formula for inclusive breakup revisited. *Eur. Phys. J. A* **57**, 57 (2021).

Transfer to continuum 模型

“There it was shown to hold under the hypothesis that the breakup process is limited to peripheral projectile-target trajectories and that it is due mainly to the neutron interaction with the target potential.”[2]

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[2] Gómez-Ramos, M., Gómez-Camacho, J., Lei, J. *et al.* The Hussein–McVoy formula for inclusive breakup revisited. *Eur. Phys. J. A* **57**, 57 (2021).